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Optical Investigations on The Transition Metal Selenides $NbSe_3$ And $(Tase_4)_2I$

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OPTICAL INVESTIGATIONS ON THE TRANSITION METAL
SELENIDES NbSe_3 AND $(\text{TaSe}_4)_2\text{I}$ *

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Abstract The polarized reflectance spectra of NbSe_3
and $(\text{TaSe}_4)_2\text{I}$ at 300 K are reported. The experimen-
tal data are analyzed in terms of a Lorentz-Drude
and a tight binding model.

INTRODUCTION

The one-dimensional metallic character of the transition metal chain compounds NbSe_3 and $(\text{TaSe}_4)_2\text{I}$ manifests itself by the formation of charge density waves which are closely connected with phase transitions and non-linear transport phenomena^{1,2}. However, in spite of this common feature the structural anisotropy of both chain compounds is quite different. Whereas NbSe_3 shows an intrachain

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metal-metal distance of 3.48 \AA and interchain metal-metal distances between 4.25 and 4.45 \AA , the corresponding values in $(\text{TaSe}_4)_2\text{I}$ are 3.21 \AA and 6.7 \AA , revealing a clearly higher anisotropy^{3,4}. In the present contribution the transport properties of both compounds along the metallic axis are compared by optical investigations.

EXPERIMENTAL RESULTS AND DISCUSSION

Fig.1 shows the polarized reflectance spectra with a metallic plasma edge for the polarization direction parallel to the metallic chains, and the nearly constant re-

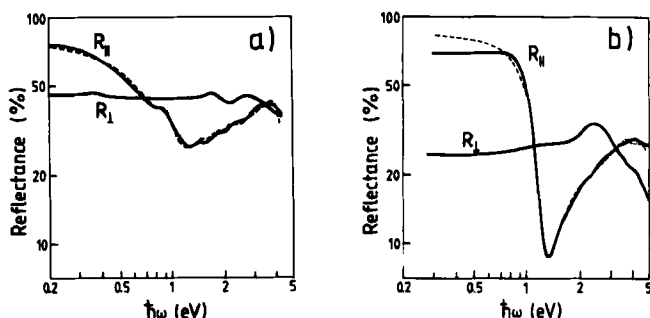


FIGURE 1 Polarized reflectance spectra of NbSe_3 (a)⁵ and $(\text{TaSe}_4)_2\text{I}$ (b). Dashed lines: Lorentz-Drude fit.

flectivity of a semiconductor for R_{\perp} . The smoother plasma edge of NbSe_3 indicates a shorter collision time of the free carriers in this material. Analyzing R_{\parallel} by a Lorentz-Drude model, the spectral distribution of the real and the imaginary part of the dielectric function, $\epsilon_1(\hbar\omega)$ and $\epsilon_2(\hbar\omega)$, are obtained (Fig.2a,b). From the relation $\sigma(\hbar\omega) = \epsilon_0 \omega \epsilon_2(\hbar\omega)$ one obtains the conductivity function, which can be extrapolated to zero energy, leading to an optical value for the dc-conductivity. As the free carrier concentration can be deduced from a discussion

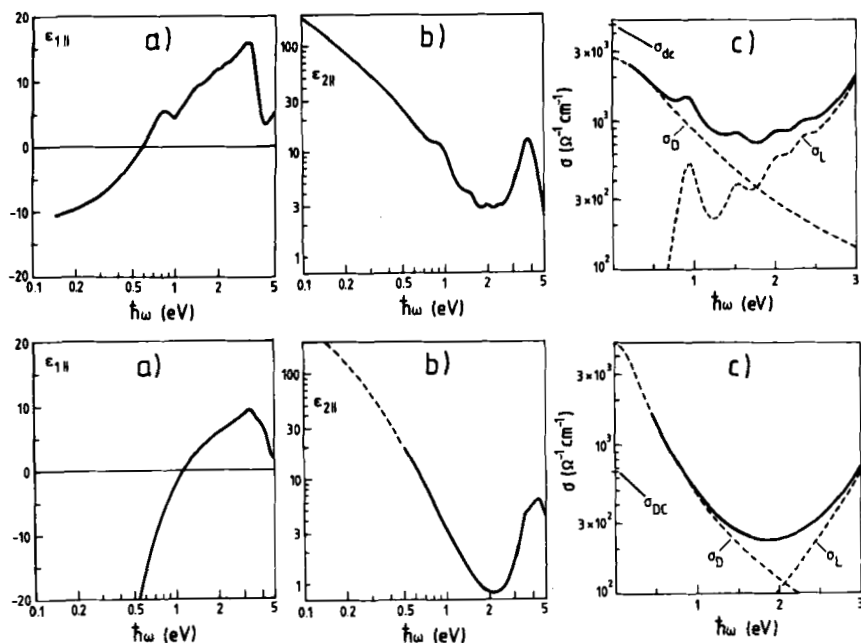


FIGURE 2 Dielectric function (a,b) and conductivity function (c) of NbSe_3 (upper part) and of $(\text{TaSe}_4)_2\text{I}$ (lower part). σ_D , σ_L and Lorentz contribution.

of the chemical bonds¹, also the effective mass can be calculated from the optical results. On the basis of a simple tight binding model follows further the band width $4t$, the Fermi energy E_F , and the mean free path Λ .

TABLE I Transport properties of NbSe_3 and $(\text{TaSe}_4)_2\text{I}$

	N ($10^{21}/\text{cm}^3$)	$\frac{m^*}{m_0}$	τ (10^{-15}sec)	$4t$ (eV)	E_F (eV)	$\frac{\Lambda}{a}$
NbSe_3	3.9	0.4	1	3.5	0.5	7
$(\text{TaSe}_4)_2\text{I}$	3.2	0.4	2	4.0	0.6	15

The most surprising results are the short mean free path of NbSe_3 and the high optical value of $4 \times 10^3 (\Omega \text{cm})^{-1}$ for the dc-conductivity of $(\text{TaSe}_4)_2\text{I}$, which is 6 times larger than the experimental dc-value.² The latter result, perhaps, indicates an influence of fluctuations, due to the phase transition at 263 K, onto the dc-conductivity at 300 K⁶. On the other hand, the value of the width of the conduction band of $(\text{TaSe}_4)_2\text{I}$ is in excellent agreement with band structure calculations if one takes into account that the tight-binding band must be folded back into a smaller Brillouin zone because the crystallographic unit cell contains four tantalum atoms along the chain axis⁷. On the contrary the width of the conduction band of NbSe_3 given here is considerably higher than the values of about 1.5 eV resulting from band-structure calculations⁸⁻¹¹. This is presumably due to the more complicated band structure of NbSe_3 with six nonequivalent metal chains.

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